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## COMMENT

# On solutions of the Yang-Baxter equations without additivity 

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#### Abstract

The relationship is described between solutions found by Ge and Xue and by the present author.


Recently two $4 \times 4$ solutions of the Yang-Baxter equations (YBE)

$$
\begin{equation*}
R_{12}(\lambda, \mu) R_{13}(\lambda, \nu) R_{23}(\mu, \nu)=R_{23}(\mu, \nu) R_{13}(\lambda, \nu) R_{12}(\lambda, \mu) \tag{1}
\end{equation*}
$$

were published [1]. The solutions are non-additive in the sense that $R(\lambda, \mu) \neq f(\lambda-\mu)$. The purpose of this comment is to generalize these solutions and describe their relationship to the solutions

$$
\begin{align*}
& R_{\mathrm{V}}(u, v)=\left(\begin{array}{cccc}
u / v & 0 & 0 & 0 \\
0 & (u v)^{-1} & 0 & 0 \\
0 & 1-k & k u v & 0 \\
0 & 0 & 0 & v / u
\end{array}\right) \quad k=\text { constant }  \tag{2}\\
& R_{\mathrm{VI}}(u, v)=\left(\begin{array}{cccc}
u_{1} / v_{2} & 0 & 0 & 0 \\
0 & \left(u_{1} v_{2}\right)^{-1} & 0 & 0 \\
0 & W & -u_{1} v_{2} & 0 \\
0 & 0 & 0 & v_{2} / u_{1}
\end{array}\right) \quad W=u_{1} / u_{2}+u_{2} / u_{1} \tag{3}
\end{align*}
$$

that are given in table 1 of [2] together with other non-additive solutions to the ybe. Note that the variables $u, v$ in $R_{\mathrm{VI}_{1}}$ are two-component quantities.

To solve the equation (1) the ansatz

$$
R(\lambda, \mu)=\left(\begin{array}{cccc}
u_{+}(\lambda, \mu) & 0 & 0 & 0  \tag{4}\\
0 & p^{(+-)}(\lambda, \mu) & 0 & 0 \\
0 & W(\lambda, \mu) & p^{(-+)}(\lambda, \mu) & 0 \\
0 & 0 & 0 & u_{-}(\lambda, \mu)
\end{array}\right)
$$

was accepted (for easy orientation we use the notation of [1]). The ansatz can be justified either [1] by weight-conservation or [2] by the requirement that we look for solutions that for $\lambda=\mu$ are of the form

$$
R=q\left(\begin{array}{cccc}
1 & 0 & 0 & 0  \tag{5}\\
0 & r & 0 & 0 \\
0 & 1-r t & t & 0 \\
0 & 0 & 0 & s
\end{array}\right) \quad s=1 \text { or } s=-r t
$$

An important fact exploited in the following is that the set of solutions of (1) is in general invariant under the transformations

$$
\begin{align*}
& R(\lambda, \mu) \mapsto \varphi(\lambda, \mu) R(\lambda, \mu)  \tag{6}\\
& R(\lambda, \mu) \mapsto[T(\lambda) \otimes T(\mu)] R(\lambda, \mu)[T(\lambda) \otimes T(\mu)]^{-1}  \tag{7}\\
& R(\lambda, \mu) \mapsto R(f(\lambda), f(\mu)) \tag{8}
\end{align*}
$$

where $\varphi$ and $f$ are scalar functions and $T$ is a $G L(2)$-valued function.
We can exploit the symmetry (6) to set $u_{+}(\lambda, \mu)=1$. Then we immediately get from (1) that $p^{(+-)}, p^{(-+)}$are functions of one variable only

$$
\begin{align*}
& p^{(+-)}(\lambda, \nu)=p^{(+-)}(\lambda, \mu)=p^{+}(\lambda)  \tag{9}\\
& p^{(-+)}(\lambda, \nu)=p^{(-+)}(\mu, \nu)=p^{-}(\nu) \tag{10}
\end{align*}
$$

(cf (11), (12) in [1]). For $W(\lambda, \mu)$ we get the equation

$$
\begin{equation*}
W(\lambda, \mu) W(\mu, \nu)=W(\lambda, \nu)\left[1-p^{+}(\mu) p^{-}(\mu)\right] \tag{11}
\end{equation*}
$$

the general solution of which is

$$
\begin{equation*}
W(\lambda, \mu)=\left[1-p^{+}(\lambda) p^{-}(\lambda)\right] \xi(\lambda) / \xi(\mu) \tag{12}
\end{equation*}
$$

where $\xi$ is an arbitrary function. The equations for $u_{-}(\lambda, \mu)$ then imply that

$$
\begin{equation*}
u_{-}(\lambda, \mu)=p^{+}(\lambda) q(\mu) \tag{13}
\end{equation*}
$$

where

$$
\begin{equation*}
q(\mu)=1 / p^{+}(\mu) \quad p^{+}(\mu) p^{-}(\mu)=k \in \mathbb{C} \backslash\{0\} \tag{14}
\end{equation*}
$$

or

$$
\begin{equation*}
q(\mu)=-p^{-}(\mu) \tag{15}
\end{equation*}
$$

The conclusion is that there are just two solutions to the ybe (1) of the form (4). They are

$$
\begin{array}{r}
R_{1}(\lambda, \mu)=\varphi(\lambda, \mu)\left(\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & p^{+}(\lambda) & 0 & 0 \\
0 & (1-k) \xi(\lambda) / \xi(\mu) & k / p^{+}(\mu) & 0 \\
0 & 0 & 0 & p^{+}(\lambda) / p^{+}(\mu)
\end{array}\right) \\
R_{2}(\lambda, \mu)=\varphi(\lambda, \mu)\left(\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & p^{+}(\lambda) & 0 & 0 \\
0 & W(\lambda, \mu) & p^{-}(\mu) & 0 \\
0 & 0 & 0 & -p^{+}(\lambda) p^{-}(\mu)
\end{array}\right) \tag{17}
\end{array}
$$

where $W$ is given by (12) and $p^{+}, p^{-}, \varphi$ and $\xi$ are arbitrary functions.
The solutios in [1] are particular cases of (16), (17) where

$$
\begin{array}{ll}
p^{+}(\lambda)=q^{-1} \eta Q^{\Sigma_{j=1}^{\prime \prime \prime} \dot{x}_{i} \lambda^{\prime}} & p^{-}(\mu)=q^{-2} / p^{+}(\mu) \\
\varphi(\lambda, \mu)=q \varphi_{+}(\lambda, \mu) & \xi(\lambda)=g(\lambda) \quad k=q^{-2} \tag{19}
\end{array}
$$

The appearance of functions $\varphi$ and $\xi$ in (16), (17) is a consequence of the symmetries (6), (7). The symmetry (8) enables to consider $p^{ \pm}(\lambda)$ and $p^{ \pm}(\mu)$ as independent variables of the solutions. This attitude was accepted in [2]. Namely, denoting $p^{+}(\lambda)=u^{-2}$, $p^{+}(\mu)=v^{-2}$ and choosing $\varphi(\lambda, \mu)=u / v, \xi(\lambda)=u^{-1}, \xi(\mu)=v^{-1}$ we get $R_{1}(\lambda, \mu)=$ $R_{\mathrm{v}}(u, v)$. Similarly, denoting $p^{+}(\lambda)=u_{1}^{-2}, p^{-}(\lambda)=-u_{2}^{2}, p^{+}(\mu)=v_{1}^{-2}, p^{-}(\mu)=-v_{2}^{2}$ and choosing $\varphi(\lambda, \mu)=u_{1} / v_{2}, \xi(\lambda)=u_{2}^{-1}, \xi(\mu)=v_{2}^{-1}$ we get $R_{2}(\lambda, \mu)=R_{\mathrm{V}_{1}}(u, v)$.

## References

[1] Ge M L and Xue K 1991 J. Phys. A: Math. Gen. 24 L895
[2] Hlavatý L 1987 J. Phys. A: Math. Gen. 201661

