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COMMENT

On solutions of the Yang-Baxter equations without additivity

Ladislav Hlavatý

Institute of Physics, Czechosłovak Academy of Sciences, Na Slovance 2, 18040 Prague 8, Czechoslovakia

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Abstract. The relationship is described between solutions found by Ge and Xue and by the present author.

Recently two 4×4 solutions of the Yang-Baxter equations (YBE)

$$R_{12}(\lambda,\mu)R_{13}(\lambda,\nu)R_{23}(\mu,\nu) = R_{23}(\mu,\nu)R_{13}(\lambda,\nu)R_{12}(\lambda,\mu)$$
(1)

were published [1]. The solutions are non-additive in the sense that $R(\lambda, \mu) \neq f(\lambda - \mu)$. The purpose of this comment is to generalize these solutions and describe their relationship to the solutions

$$R_{\rm v}(u,v) = \begin{pmatrix} u/v & 0 & 0 & 0 \\ 0 & (uv)^{-1} & 0 & 0 \\ 0 & 1-k & kuv & 0 \\ 0 & 0 & 0 & v/u \end{pmatrix} \qquad k = \text{constant} \qquad (2)$$

$$R_{\rm vI}(u,v) = \begin{pmatrix} u_1/v_2 & 0 & 0 & 0 \\ 0 & (u_1v_2)^{-1} & 0 & 0 \\ 0 & W & -u_1v_2 & 0 \\ 0 & 0 & 0 & v_2/u_1 \end{pmatrix} \qquad W = u_1/u_2 + u_2/u_1 \qquad (3)$$

that are given in table 1 of [2] together with other non-additive solutions to the YBE. Note that the variables u, v in R_{VI} are two-component quantities.

To solve the equation (1) the ansatz

$$R(\lambda,\mu) = \begin{pmatrix} u_{+}(\lambda,\mu) & 0 & 0 & 0\\ 0 & p^{(+-)}(\lambda,\mu) & 0 & 0\\ 0 & W(\lambda,\mu) & p^{(-+)}(\lambda,\mu) & 0\\ 0 & 0 & 0 & u_{-}(\lambda,\mu) \end{pmatrix}$$
(4)

was accepted (for easy orientation we use the notation of [1]). The ansatz can be justified either [1] by weight-conservation or [2] by the requirement that we look for solutions that for $\lambda = \mu$ are of the form

$$R = q \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & r & 0 & 0 \\ 0 & 1 - rt & t & 0 \\ 0 & 0 & 0 & s \end{pmatrix} \qquad s = 1 \text{ or } s = -rt.$$
(5)

An important fact exploited in the following is that the set of solutions of (1) is in general invariant under the transformations

$$R(\lambda,\mu) \mapsto \varphi(\lambda,\mu) R(\lambda,\mu) \tag{6}$$

$$R(\lambda,\mu) \mapsto [T(\lambda) \otimes T(\mu)] R(\lambda,\mu) [T(\lambda) \otimes T(\mu)]^{-1}$$
(7)

$$R(\lambda,\mu) \mapsto R(f(\lambda), f(\mu)) \tag{8}$$

where φ and f are scalar functions and T is a GL(2)-valued function.

We can exploit the symmetry (6) to set $u_+(\lambda, \mu) = 1$. Then we immediately get from (1) that $p^{(+-)}$, $p^{(-+)}$ are functions of one variable only

$$p^{(+-)}(\lambda, \nu) = p^{(+-)}(\lambda, \mu) = p^{+}(\lambda)$$
(9)

$$p^{(-+)}(\lambda, \nu) = p^{(-+)}(\mu, \nu) = p^{-}(\nu).$$
(10)

(cf (11), (12) in [1]). For $W(\lambda, \mu)$ we get the equation

$$W(\lambda, \mu) W(\mu, \nu) = W(\lambda, \nu) [1 - p^{+}(\mu) p^{-}(\mu)]$$
(11)

the general solution of which is

$$W(\lambda,\mu) = [1 - p^{+}(\lambda)p^{-}(\lambda)]\xi(\lambda)/\xi(\mu)$$
(12)

where ξ is an arbitrary function. The equations for $u_{-}(\lambda, \mu)$ then imply that

$$u_{-}(\lambda,\mu) = p^{+}(\lambda)q(\mu) \tag{13}$$

where

$$q(\mu) = 1/p^{+}(\mu)$$
 $p^{-}(\mu) p^{-}(\mu) = k \in \mathbb{C} \setminus \{0\}$ (14)

ог

$$q(\mu) = -p^{-}(\mu). \tag{15}$$

The conclusion is that there are just two solutions to the YBE(1) of the form (4). They are

$$R_{1}(\lambda,\mu) = \varphi(\lambda,\mu) \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & p^{+}(\lambda) & 0 & 0 \\ 0 & (1-k)\xi(\lambda)/\xi(\mu) & k/p^{+}(\mu) & 0 \\ 0 & 0 & 0 & p^{+}(\lambda)/p^{+}(\mu) \end{pmatrix}$$
(16)
$$R_{2}(\lambda,\mu) = \varphi(\lambda,\mu) \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & p^{+}(\lambda) & 0 & 0 \\ 0 & W(\lambda,\mu) & p^{-}(\mu) & 0 \\ 0 & 0 & 0 & -p^{+}(\lambda)p^{-}(\mu) \end{pmatrix}$$
(17)

where W is given by (12) and p^+ , p^- , φ and ξ are arbitrary functions.

The solutios in [1] are particular cases of (16), (17) where

$$p^{+}(\lambda) = q^{-1} \eta Q^{\sum_{j=1}^{m} \hat{a}_{j} \lambda^{j}} \qquad p^{-}(\mu) = q^{-2} / p^{+}(\mu).$$
(18)

$$\varphi(\lambda,\mu) = q\varphi_+(\lambda,\mu)$$
 $\xi(\lambda) = g(\lambda)$ $k = q^{-2}$ (19)

The appearance of functions φ and ξ in (16), (17) is a consequence of the symmetries (6), (7). The symmetry (8) enables to consider $p^{\pm}(\lambda)$ and $p^{\pm}(\mu)$ as independent variables of the solutions. This attitude was accepted in [2]. Namely, denoting $p^{+}(\lambda) = u^{-2}$, $p^{+}(\mu) = v^{-2}$ and choosing $\varphi(\lambda, \mu) = u/v$, $\xi(\lambda) = u^{-1}$, $\xi(\mu) = v^{-1}$ we get $R_1(\lambda, \mu) =$ $R_{\nu}(u, v)$. Similarly, denoting $p^{+}(\lambda) = u_1^{-2}$, $p^{-}(\lambda) = -u_2^2$, $p^{+}(\mu) = v_1^{-2}$, $p^{-}(\mu) = -v_2^2$ and choosing $\varphi(\lambda, \mu) = u_1/v_2$, $\xi(\lambda) = u_2^{-1}$, $\xi(\mu) = v_2^{-1}$ we get $R_2(\lambda, \mu) = R_{\nu 1}(u, v)$.

References

- [1] Ge M L and Xue K 1991 J. Phys. A: Math. Gen. 24 L895
- [2] Hlavatý L 1987 J. Phys. A: Math. Gen. 20 1661